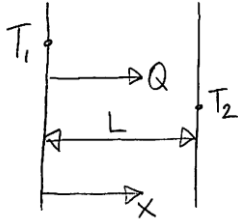


## One-Dimensional, Steady-State Heat Conduction (Reorganization of the Lecture Notes from Professor Nenad Miljkovic)

### 1-D, steady state, $\dot{Q}''' = 0$ , $k = \text{constant}$

We know from heat diffusion equation that  $\nabla^2 T = 0$ .

(1) Slab



$$\int \frac{\partial^2 T}{\partial x^2} = \int 0$$

$$\int \frac{\partial T}{\partial x} = \int C_1$$

$$T(x) = C_1 x + C_2$$

Given the boundary conditions  $T(x=0) = T_1$ ,  $T(x=L) = T_2$

$$T(x) = \frac{T_2 - T_1}{L} * x + T_1$$

So if we define a heat transfer resistance:  $R$ ,  $A = \text{area}$

$$Q = \frac{\Delta T}{R}$$

$$Q = -kA \frac{\partial T}{\partial x}$$

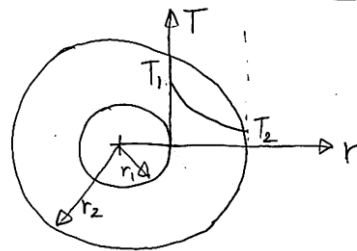
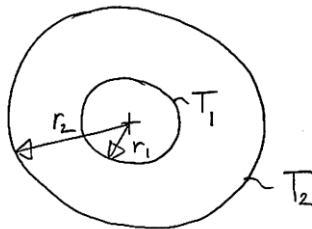
From above we know that  $\frac{\partial T}{\partial x} = \frac{T_2 - T_1}{L} = -\frac{\Delta T}{L}$

$$Q = \frac{kA}{L} * \Delta T = \frac{1}{R} * \Delta T$$

Thermal resistance of a solid slab:

$$R_{\text{slab}} = L/kA$$

(2) Circular cylinder



We know the heat equation in radial coordinate is:

$$\frac{1}{r} * \frac{\partial}{\partial r} \left( r * \frac{\partial T}{\partial r} \right) = 0$$

$$r \frac{\partial T}{\partial r} = C_1$$

$$T(r) = C_1 \ln(r) + C_2$$

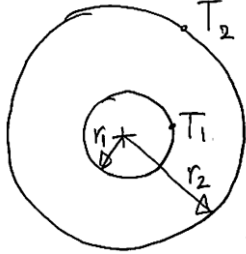
After applying the boundary conditions:

$$\frac{T - T_1}{T_2 - T_1} = \ln\left(\frac{r}{r_1}\right) / \ln\left(\frac{r_2}{r_1}\right)$$

Similarly, cylindrical thermal resistance:

$$R_{cyl} = |\ln(r_2/r_1) / 2\pi kL|$$

(3) Spherical system



Spherical heat diffusion equation:

$$\frac{1}{r^2} * \frac{\partial (r^2 * \frac{\partial T}{\partial r})}{\partial r} = 0$$

Spherical temperature profile:

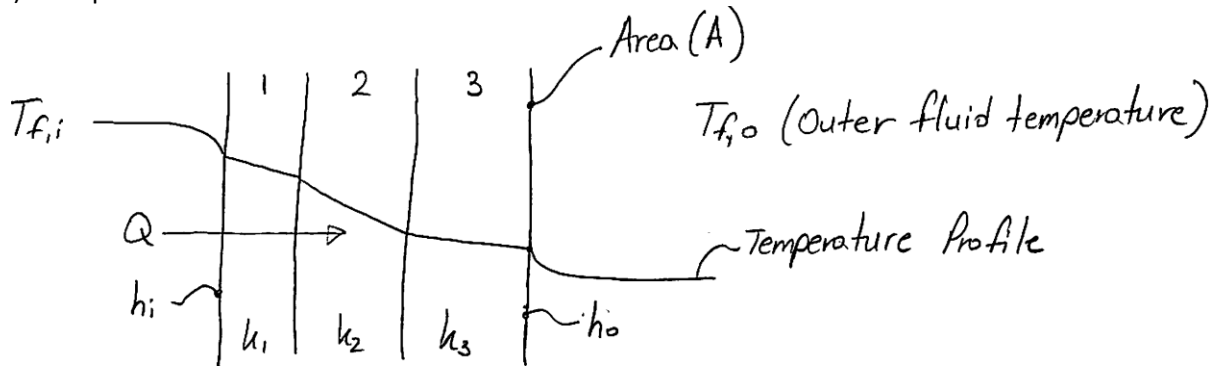
$$\frac{T - T_2}{T_1 - T_2} = \left(\frac{1}{r} - \frac{1}{r_2}\right) / \left(\frac{1}{r_1} - \frac{1}{r_2}\right)$$

Spherical thermal resistance:

$$R_{sph} = \frac{|\frac{1}{r_1} - \frac{1}{r_2}|}{4\pi k}$$

**Composite problems**

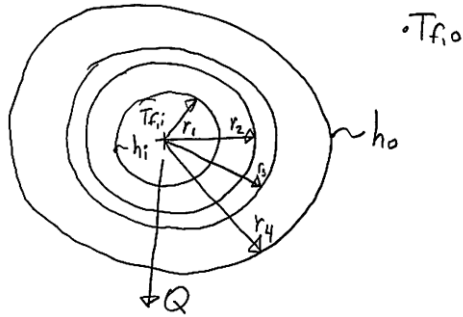
(1) Composite wall



$$Q = \frac{A(T_{f,i} - T_{f,o})}{\frac{1}{h_i} + \frac{1}{h_o} + \sum_{i=1}^n \frac{L_i}{k_i}}$$

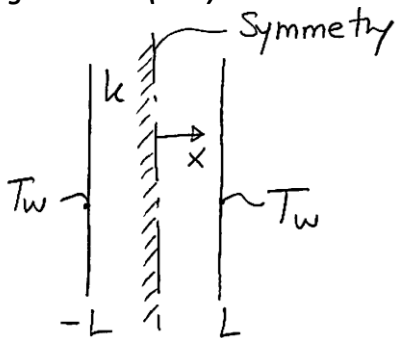
- $L_i$ =thickness of the i-th wall
- $k_i$ =thermal conductivity of the i-th wall
- $h_i$ =inner wall heat transfer coefficient
- $h_o$ =outer wall heat transfer coefficient

(2) Composite cylinder



$$Q = \frac{2\pi l(T_{f,i} - T_{f,o})}{\frac{1}{h_i r_i} + \frac{1}{h_o r_{n+1}} + \sum_{j=1}^n \frac{\ln\left(\frac{r_{j+1}}{r_j}\right)}{k_j}}$$

**Heat generation (slab)**



Steady state, 1-D, constant properties, uniform  $\dot{Q}'''$ .

$$\frac{\partial(k * \frac{\partial T}{\partial x})}{\partial x} + \dot{Q}''' = \frac{\rho C_p \partial T}{\partial t} = 0$$

$$\frac{\partial^2 T}{\partial x^2} + \frac{\dot{Q}'''}{k} = 0$$

Boundary conditions:

$$T(x=L) = T_w$$

$$\frac{\partial T}{\partial x} |_{x=0} = 0$$

$$T(x) = T_w + \frac{\dot{Q}''' L^2}{2k} \left(1 - \left(\frac{x}{L}\right)^2\right)$$

$$T_{\max} = T_w + \frac{\dot{Q}''' L^2}{2k} \text{ at } x=0.$$

$$q''_{\text{out}} = \dot{Q}''' L$$

In general, for heat generation we have the following formulation:

$$-\frac{k}{r^n} \frac{\partial(r^n * \frac{\partial T}{\partial r})}{\partial r} = \dot{Q}'''$$

Where  $n=0$  for slab,  $n=1$  for cylinder,  $n=2$  for sphere.

This only works if  $\dot{Q}''' = \text{constant}$

Our generalized result is:

Temperature profile:

$$T(r) = T_a + \frac{\dot{Q}''' a^2}{2(n+1)k} \left(1 - \frac{r^2}{a^2}\right)$$

Heat flux:

$$q'' = \frac{\dot{Q}''' r}{n + 1}$$

Maximum temperature where  $T_a$  is the boundary temperature on the outside:

$$T_{\max} = T(r = 0) = T_a + \frac{\dot{Q}''' a^2}{2(n + 1)k}$$

**Conduction for general shapes**

No convection, steady state,  $\dot{Q}''' = 0$ ,  $k = \text{constant}$ , 1-D

Define an orthogonal coordinate system:

$u_1, u_2, u_3$  (directions)  $s_1, s_2, s_3$  (lengths)  $ds_1 = h_1 du_1, ds_2 = h_2 du_2, ds_3 = h_3 du_3$

$$q'' = -k * \frac{\partial T}{\partial s_1}$$

$$Q = Aq'' = -kA(u_1) \frac{\partial T}{\partial s_1}$$

Rearranging and integrating (assuming  $ds_1 = du_1$ )

$$\frac{Q}{k} * \int_{(u_1)_a}^{(u_1)_b} \frac{du_1}{A(u_1)} = - \int_{T_a}^{T_b} dT = T_a - T_b = \Delta T$$

$$Q = \Delta T / \left( \frac{1}{k} * \int \frac{du_1}{A(u_1)} \right) = \Delta T / R$$

$$R = \frac{1}{k} * \int \frac{du_1}{A(u_1)}$$

Double check our previous solutions:

(1) Slab:  $u_1 = x, du_1 = dx, A(u_1) = \text{constant}$

$$R_{\text{slab}} = \frac{1}{kA} \int du_1 = \frac{1}{kA} \int dx = L/kA$$

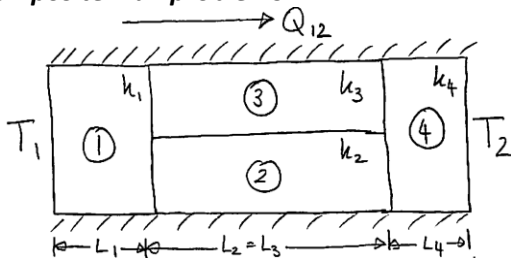
(2) Cylinder:  $u_1 = r, du_1 = dr, A(u_1) = A(r) = 2\pi r l$

$$R_{\text{cyl}} = \frac{1}{k} * \int_{r_1}^{r_2} \frac{dr}{2\pi r l} = \frac{\ln\left(\frac{r_2}{r_1}\right)}{2\pi k l}$$

(3) Sphere:  $u_1 = r, du_1 = dr, A(u_1) = A(r) = 4\pi r^2$

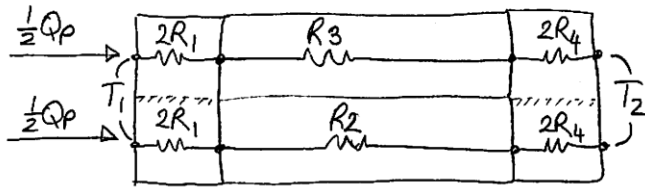
$$R_{\text{sph}} = \frac{1}{k} * \int_{r_1}^{r_2} \frac{dr}{4\pi r^2} = \frac{1}{4\pi k} * \int_{r_1}^{r_2} \frac{dr}{r^2} = \frac{1}{4\pi k} (1/r_1 - 1/r_2)$$

**Composite wall problems**



There are two ways to analyze this problem. They give different answers.

(1) Parallel approach



Add adiabatic surface to divide  $R_1$  and  $R_4$  into two.

$$R_{\text{Lower}} = \frac{L_1}{k_1 A_2} + \frac{L_2}{k_2 A_2} + \frac{L_4}{k_4 A_2}$$

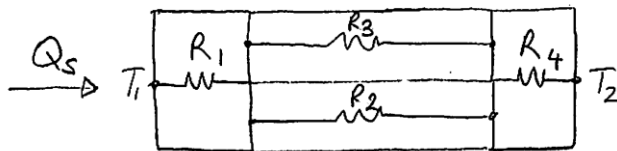
$$R_{\text{Upper}} = \frac{L_1}{k_1 A_3} + \frac{L_2}{k_3 A_3} + \frac{L_4}{k_4 A_3}$$

$$\frac{1}{R_p} = \frac{1}{R_{\text{Lower}}} + \frac{1}{R_{\text{Upper}}}$$

$$\frac{1}{R_p} = \left( \frac{A_2}{\frac{L_1}{k_1} + \frac{L_2}{k_2} + \frac{L_4}{k_4}} + \frac{A_3}{\frac{L_1}{k_1} + \frac{L_2}{k_3} + \frac{L_4}{k_4}} \right)$$

$$Q_p = \frac{\Delta T}{R_p} = \left( \frac{A_2}{\frac{L_1}{k_1} + \frac{L_2}{k_2} + \frac{L_4}{k_4}} + \frac{A_3}{\frac{L_1}{k_1} + \frac{L_2}{k_3} + \frac{L_4}{k_4}} \right) \Delta T$$

(2) Series approach



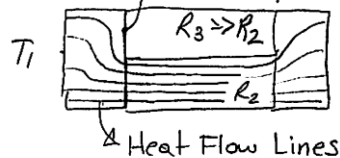
$$R_s = \frac{L_1}{k_1 A_1} + \left( \frac{k_2 A_2}{L_2} + \frac{k_3 A_3}{L_2} \right)^{-1} + \frac{L_4}{k_4 A_1}$$

$$Q_s = \frac{\Delta T}{R_s} = \left( \frac{L_1}{k_1 A_1} + \left( \frac{k_2 A_2}{L_2} + \frac{k_3 A_3}{L_2} \right)^{-1} + \frac{L_4}{k_4 A_1} \right)^{-1} \Delta T$$

It is always true that  $Q_p \leq Q_{\text{real}} \leq Q_s$

Aside: Think of this way:

Real Life:  $T$  will vary here



While by adding infinite conductance between  $R_1$  and  $R_2$ ,  $R_3$ , the temperature at the boundary becomes the same.

In general:

1. Adding adiabatic surface:  $R$  goes up,  $Q$  goes down;
2. Allowing infinite conductance:  $R$  goes down,  $Q$  goes up.